

## Reference crust-mantle density contrast beneath Antarctica based on the Vening Meinesz-Moritz isostatic inverse problem and CRUST2.0 seismic model

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### ABSTRACT

The crust-mantle (Moho) density contrast beneath Antarctica was estimated based on solving the Vening Meinesz-Moritz isostatic problem and using constraining information from a seismic global crustal model (CRUST2.0). The solution was found by applying a least-squares adjustment by elements method. Global geopotential model (GOCO02S), global topographic/bathymetric model (DTM2006.0), ice-thickness data for Antarctica (assembled by the BEDMAP project) and global crustal model (CRUST2.0) were used for computing isostatic gravity anomalies. Since CRUST2.0 data for crustal structures under Antarctica are not accurate (due to a lack of seismic data in this part of the world), Moho density contrast was determined relative to a reference homogenous crustal model having  $2,670 \text{ kg/m}^3$  constant density. Estimated values of Moho density contrast were between  $160$  and  $682 \text{ kg/m}^3$ . The spatial distribution of Moho density contrast resembled major features of the Antarctic's continental and surrounding oceanic tectonic plate configuration; maxima exceeding  $500 \text{ kg/m}^3$  were found throughout the central part of East Antarctica, with an extension beneath the Transantarctic mountain range. Moho density contrast in West Antarctica decreased to  $400\text{-}500 \text{ kg/m}^3$ , except for local maxima up to  $\sim 550 \text{ kg/m}^3$  in the central Antarctic Peninsula.

*Key words:* Antarctica, crust, gravity, isostasy, mantle, Moho interface.

### RESUMEN

El contraste de densidad de la discontinuidad de Mohorovičić (Moho) debajo de la Antártida fue estimado con base en la solución del problema isostático Vening Meinesz-Moritz y a partir de datos obtenidos con el modelo sísmico de la corteza global (CRUST2.0). La solución se encontró a través de un ajuste al método de mínimos cuadrados por el método de elementos. El modelo geopotencial global (GOCO02S), el modelo topográfico/batimétrico (DTM2006.0), los datos de espesor del hielo para la Antártida (reunidos por el proyecto BEDMAP) y el modelo sísmico de corteza global (CRUST2.0) fueron utilizados para calcular las anomalías gravitatorias isostáticas. Ya que los datos de CRUST2.0 para las estructuras de la corteza en la Antártida no son exactos (debido a la falta de información sísmica para esta parte del planeta), el contraste de densidad de la Discontinuidad de Mohorovičić fue determinado a partir de un modelo de corteza homogéneo que tiene una densidad constante de  $2,670 \text{ kg/m}^3$ . Los valores estimados del contraste de densidad de la Moho se encontraron entre  $160$  y  $682 \text{ kg/m}^3$ . La distribución espacial del contraste de densidad de la Moho exhibe mayores rasgos en la configuración de la plancha tectónica de la Antártida continental y su alrededor oceánico. El valor máximo encontrado excede los  $500 \text{ kg/m}^3$  y se ubica en la parte Este continental, con extensión en las Montañas Transantárticas. El contraste de densidad de la Moho (zona de transición entre la corteza y el manto terrestre) en el Oeste de la Antártida osciló entre  $400\text{-}500 \text{ kg/m}^3$ , excepto para la máxima local de  $\sim 550 \text{ kg/m}^3$ , en el centro de la Península Antártida.

*Palabras clave:* Antártida, corteza terrestre, manto terrestre, gravedad, isostasia, Discontinuidad de Mohorovičić

*Record*

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### Introduction

Current knowledge about the lithospheric structure beneath Antarctica is limited due to a low spatial coverage of high-quality seismic data. Seismic studies by Kogan (1972) and Ito and Ikami (1986) were

based on localised controlled source seismic experiments. Passive seismic studies, based on earthquakes occurring mostly outside the Antarctic tectonic plate (due to a lack of intra-plate seismicity within the Antarctic plate (Okal, 1981), still represent the primary source of information. Studies based on an analysis of surface wave velocity can be found in

Evison *et al.*, (1960), Kovach and Press (1961), Bentley and Ostenso (1962), Dewart and Toksoz (1965), Adams (1971), Knopoff and Vane (1978), Rouland *et al.*, (1985), Forsyth *et al.*, (1987), Roullet *et al.*, (1994) and Bannister *et al.*, (2003). Seismic receiver function analysis has been carried out by Reading (2006), Lawrence *et al.*, (2006) and Winberry and Anandakrishnan (2004). Ritzwoller *et al.*, (2001) used the simultaneous inversion of broadband group velocity measurements to compile a seismic model of the crust and upper mantle beneath Antarctica and surrounding oceans. Some studies have been based on airborne gravity surveys, for instance by Studinger *et al.*, (2004, 2006). Llubes *et al.*, (2003) used CHAMP satellite gravity data for estimating crust thickness in Antarctica. More recently, Block *et al.*, (2009) have estimated crust thickness in Antarctica using GRACE gravity data.

Despite some authors having investigated Antarctic crust thickness, studies addressing lithosphere density structure and density interface in this part of the world are rare. This study was thus aimed at investigating Moho density contrast beneath Antarctic continental and surrounding oceanic crustal structures. Since large areas of Antarctica are not yet sufficiently covered by seismic surveys, current knowledge about crust density structure is not complete and accurate. A simple model of 2,670 kg/m<sup>3</sup> homogenous crust reference density was thus adopted and Moho density contrast was estimated regarding such crustal density. This involved applying the method recently developed by Sjöberg (2009) and Sjöberg and Bagherbandi (2011).

### Isostatic model

The generic expression of solving the Vening Meinesz-Moritz isostatic problem was formulated in the following form (Sjöberg and Bagherbandi, 2011):

$$-G R \iint_{\Phi} \Delta\rho(\Omega') K(\psi, s) d\Omega' = \Delta\tilde{g}^i(r, \Omega), \quad (1)$$

where  $G = 6.674 \times 10^{-11} \text{ m}^3 \text{ kg}^{-1} \text{ s}^{-2}$  is Newton's gravitational constant,  $R = 6371 \times 10^3 \text{ m}$  Earth's mean radius,  $\Delta\rho$  Moho density contrast,  $\Delta\tilde{g}^i$  (approximate) isostatic gravity anomaly (cf. Sjöberg, 2009),  $K$  the integral kernel function, and  $d\Omega' = \cos\phi' d\phi' d\lambda'$  the infinitesimal surface element on the unit sphere. The 3-D position was defined in the spherical coordinates system  $(r, \Omega)$ , where  $r$  is the spherical radius and  $\Omega = (\phi, \lambda)$  denotes the spherical direction with spherical latitude  $\phi$  and longitude  $\lambda$ . The full spatial angle is denoted as  $\Phi = \{\Omega' = (\phi', \lambda'): \phi' \in [-\pi/2, \pi/2] \wedge \lambda' \in [0, 2\pi]\}$ .

Integral kernel  $K$  in equation (1) was defined for parameters  $\psi$  and  $s$ , where  $\psi$  is the spherical distance between observation and (running) integration points  $(r, \Omega)$  and  $(r', \Omega')$ , and  $s$  a ratio function of the Moho depth  $T$  and Earth's mean radius  $R$ , i.e.  $s = 1 - \tau = 1 - T/R$ . The spectral representation of  $K$  was given by (Sjöberg, 2009):

$$K(\psi, s) = \sum_{n=0}^{\infty} \frac{n+1}{n+3} (1-s^{n+3}) P_n(\cos\psi), \quad (2)$$

where  $P_n$  is the Legendre polynomial of degree  $n$  for the argument of cosine of the spherical distance  $\psi$ .

The isostatic gravity anomaly  $\Delta\tilde{g}^i$  computed at position  $(r, \Omega)$  was defined as follows (cf. Vening Meinesz, 1931):

$$\Delta\tilde{g}^i(r, \Omega) = \Delta g^B(r, \Omega) + g^c(r, \Omega) = 0, \quad (3)$$

where  $\Delta g^B$  is the refined Bouguer gravity anomaly, and  $g^c$  the gravitational attraction of isostatic compensation masses (see also Bjerhammar, 1962, Chap. 14, eqn. 5; Moritz, 1990). The spectral representation of isostatic gravity anomaly  $\Delta\tilde{g}^i$  was defined in terms of spherical harmonics  $\Delta\tilde{g}_{n,m}^i$  of degree  $n$  and order  $m$ . The numerical coefficients  $\Delta\tilde{g}_{n,m}^i$  were generated from spherical harmonic coefficients of gravity anomaly  $\Delta g_{n,m}$  after applying

the Bouguer gravity reduction term  $2\pi G(\bar{\rho}^c H)_{n,m}$ , which was defined by the coefficients of global topographic/bathymetric (density) spherical functions  $(\bar{\rho}^c H)_{n,m}$ . The coefficients  $\Delta\tilde{g}_{n,m}^i$  were defined as (Sjöberg, 2009):

$$\Delta\tilde{g}_{n,m}^i = \frac{1}{4\pi} \begin{cases} 2\pi G(\bar{\rho}^c H)_{0,0} - \tilde{g}_0^c & \text{if } n=0 \\ 2\pi G(\bar{\rho}^c H)_{n,m} - \Delta g_{n,m} & \text{otherwise} \end{cases}, \quad (4)$$

where  $Y_{n,m}$  is the (fully-normalised) surface spherical harmonic function of degree  $n$  and order  $m$ . The density distribution function  $\bar{\rho}^c$  equals  $\bar{\rho}^c = \rho^c$  on land and ocean density contrast is defined as:  $\bar{\rho}^c = \rho^c - \rho^w$ ; where  $\rho^c$  is waverference crust density and  $\rho^w$  mean saltwater density. Ice and sediment density contrasts were defined relative to reference crust density  $\rho^c$ . Nominal (zero-degree) compensation attraction  $\tilde{g}_0^c$  stipulated at the sphere of radius  $R$  was computed as from (Sjöberg, 2009):

$$\begin{aligned} \tilde{g}_0^c &= g_0^c(r, \Omega) \Big|_{r=R} = \frac{4}{3} \pi G \Delta\rho_0 R [(1-\tau_0)^3 - 1] \\ &\approx -4\pi G \Delta\rho_0 T_0, \end{aligned} \quad (5)$$

where  $\tau_0 = T_0/R$ , and  $T_0$  and  $\Delta\rho_0$  are the adopted nominal mean values of Moho depth and density contrast, respectively.

### Estimation principle

Least-squares analysis was used for estimating  $T$  and  $\Delta\rho$ . The linearised observation equation involved expanding the integral term on the left-hand side of equation (1) into a Taylor series. The subsequent substitution of the first two terms in the series for  $s^{n+3}$  from equation (3) to equation (1) yielded (Sjöberg and Bagherbandi, 2011):

$$\begin{aligned} R \iint_{\Phi} \Delta\rho(\Omega') K(\psi, s) d\Omega' &= R \sum_{n=0}^{\infty} (n+1) \iint_{\Phi} \Delta\rho(\Omega') \left( \tau - \tau^2 \frac{n+2}{2} \right) P_n(\cos\psi) d\Omega' \\ &= 4\pi \sum_{n=0}^{\infty} \sum_{m=-n}^n \frac{n+1}{2n+1} \left[ (\Delta\rho T)_{n,m} - \frac{n+2}{2R} (\Delta\rho T^2)_{n,m} \right] Y_{n,m}(\Omega). \end{aligned} \quad (6)$$

From equation (5), the linearised observation equation for product  $T\Delta\rho$  was found in the following form:

$$T(\Omega)\Delta\rho(\Omega) = \sum_{n=0}^{\infty} \sum_{m=-n}^n \left[ \frac{2n+1}{4\pi G(n+1)} \Delta\tilde{g}_{n,m}^i - \frac{n+2}{2} (\Delta\rho T^2)_{n,m} \right] Y_{n,m}(\Omega). \quad (7)$$

Approximation of term  $(\Delta\rho T^2)_{n,m}$  in equation (7) by  $T_0(\Delta\rho T)_{n,m}$  yielded the linearised observation equation for  $\Delta\rho$  in the following form:

$$\Delta\rho(\Omega) \approx \frac{\Delta\tilde{g}^i(r, \Omega)}{2\pi G T(\Omega)} - \frac{1}{4\pi T(\Omega)} \sum_{n=0}^{\infty} \sum_{m=-n}^n \left[ \frac{1}{n+1} - \frac{T_0/R}{2/(n+2) - T_0/R} \right] \Delta\tilde{g}_{n,m}^i Y_{n,m}(\Omega). \quad (8)$$

Least-squares analysis combined the estimated product of  $T$  and  $\Delta\rho$  with a priori values  $t$  and  $\kappa$  of such parameters for obtaining improved estimates of  $T$  and  $\Delta\rho$ . The system of observation equations was given in the following vector-matrix form:

$$\mathbf{A} \mathbf{x} = \mathbf{l} - \boldsymbol{\varepsilon}, \quad (9)$$

where  $\boldsymbol{\varepsilon}$  is the vector of residuals. Parameter vector  $\mathbf{x}$  and observation vector  $\mathbf{l}$  read:

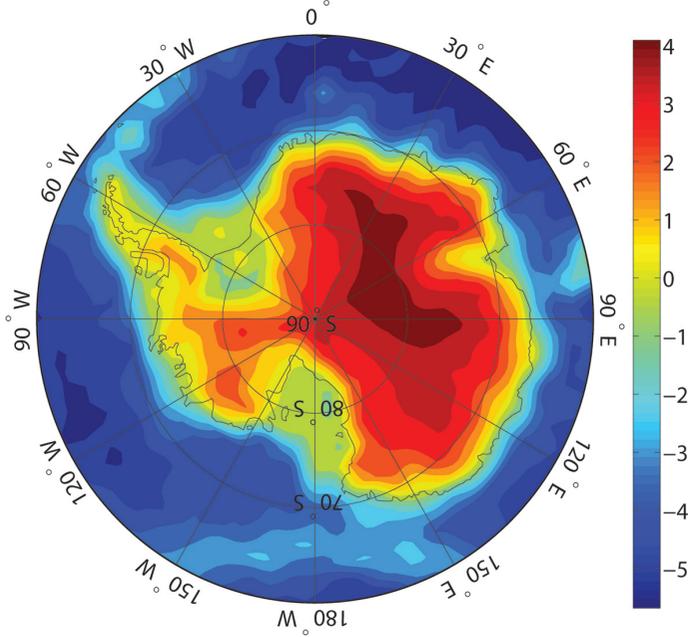
$$\mathbf{A} = \begin{pmatrix} \kappa & t \\ 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \mathbf{x} = \begin{pmatrix} dT \\ d\kappa \end{pmatrix}, \quad \mathbf{l} = \begin{pmatrix} l_1 - t\kappa \\ l_2 - t \\ l_3 - t \end{pmatrix}, \quad (10)$$

Elements  $l_1$ ,  $l_2$  and  $l_3$  of the observation vector  $\mathbf{l}$  were formed, respectively, by observables  $T\Delta\rho$ ,  $\Delta\rho$  and  $T$ . Parameter vector  $\mathbf{x}$  was defined for unknown corrections  $dT$  and  $d\kappa$  to *a priori* (initial) values of  $T$  and  $\Delta\rho$ . The solution of normal equations was found by solving:

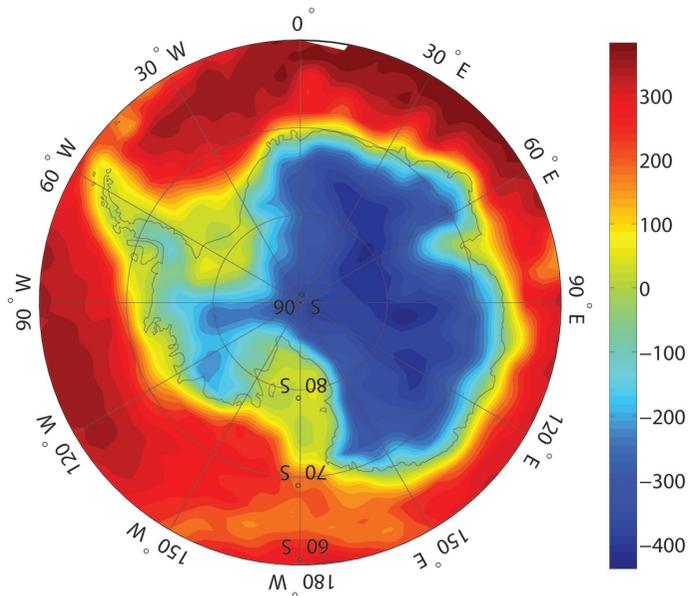
$$\hat{\mathbf{x}} = \mathbf{N}^{-1} \mathbf{A}^T \mathbf{Q}^{-1} \mathbf{l}, \mathbf{N} = \mathbf{A}^T \mathbf{Q}^{-1} \mathbf{A}, \quad (11)$$

Covariance matrix  $\mathbf{Q}_{\hat{\mathbf{x}}}$  of the estimated parameters was computed from:

$$\mathbf{Q}_{\hat{\mathbf{x}}} = \sigma_0^2 \mathbf{N}^{-1}, \quad (12)$$



**Fig. 1** The Earth's solid topography generated on a  $2 \times 2$  arc-deg grid using DTM2006.0 coefficients complete to degree/order of 180. Values are in km.



**Fig. 2** Refined Bouguer gravity anomalies computed on a  $2 \times 2$  arc-deg surface grid with a spectral resolution complete to degree 180 of spherical harmonics. Values are in miligals.

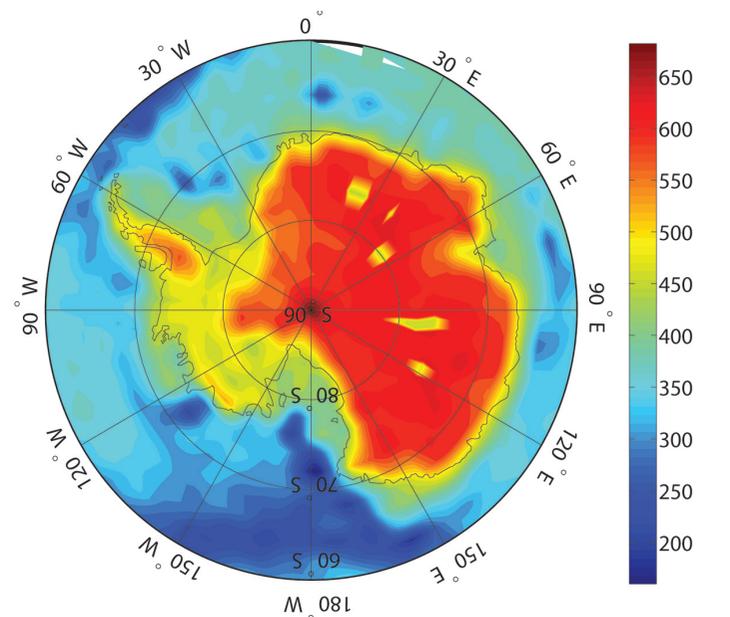
where unit weight variance  $\sigma_0^2$  reads was:

$$\sigma_0^2 = \mathbf{l}^T \mathbf{Q}^{-1} (\mathbf{l} - \mathbf{A} \hat{\mathbf{x}}). \quad (13)$$

## Data acquisition

The computation of isostatic gravity anomalies over the study area of Antarctica required the application of gravity corrections due to rough topography, continental ice sheet and variable geological structure (mainly large sediment deposits). The three largest mountain ranges on the Antarctic continent are the Transantarctic Mountains, the West and East Antarctica ranges. The Transantarctic Mountains is formed by a mountain range extending, with some interruptions, across the continent from Cape Adare in northern Victoria Land to Coats Land. These mountains serve as the natural geographical division between East and West Antarctica. Absolute maxima of topographic gravity correction in this part of the world reach  $\sim 400$  mGal (Tenzer *et al.*, 2009). Approximately 98% of the land in Antarctica is covered by the continental ice sheet, thickness exceeding 2 km, while only about 0.32% of the land is ice-free. Maximum ice thickness is 4,776 m. The ice mass thus represents a significant contribution to the gravity field. Another significant gravity field contribution along the continental shelf of Antarctica is due to the presence of large sedimentary basins. Additional large sedimentary deposits exist inland of Antarctica (see Studinger *et al.*, 2003; Bamber *et al.*, 2006). Tenzer *et al.*, (2010, 2012) have demonstrated that ice and sediment stripping corrections to gravity data in Antarctica reach about 300 mGal and 60 mGal, respectively.

The global geopotential model (GOCO02S), the global topographic/bathymetric model (DTM2006.0), ice-thickness data for Antarctica (assembled by the BEDMAP project; Lythe *et al.*, 2001) and sediment density and thickness data from the global crustal model (CRUST2.0) were used in this study to compute isostatic gravity anomalies over the study area of Antarctica bounded by parallel 60 arc-deg of southern geographical latitude. All gravity computations were realized on a  $2 \times 2$  arc-deg surface grid. Coefficients from the combined GRACE and GOCE satellite global geopotential model GOCO02S (Goiginger *et al.*, 2011) with a spectral resolution complete to degree 180 spherical harmonics were used for producing gravity anomalies. Refined Bouguer gravity anomalies were ob-



**Fig. 3** Crust-mantle density contrast computed on a  $2 \times 2$  arc-deg grid using a combined least-squares method. Values are in  $\text{kg}/\text{m}^3$ .

tained after applying Bouguer gravity reductions to GOCO02S gravity anomalies. The spherical Bouguer gravity reduction was computed using coefficients from the global topographic/bathymetric model DTM2006.0 (Pavlis *et al.*, 2007) complete to spherical harmonic degree 180. The average density of upper continental crust  $2,670 \text{ kg/m}^3$  (Hinze, 2003) was adopted as reference crust density. The ocean density contrast  $1,643 \text{ kg/m}^3$  corresponds to mean seawater density  $1,027 \text{ kg/m}^3$ . Updated  $5 \times 5$  arc-min ice-thickness data for Antarctica assembled by the BEDMAP project were used to compute the ice (density contrast) stripping gravity correction. The ice stripping gravity correction was computed with a spectral resolution complete to degree/order 180. The ice density contrast  $1,753 \text{ kg/m}^3$  corresponds to glacial ice density  $2,670 \text{ kg/m}^3$  (Cutnell and Kenneth, 1995). CRUST2.0 sediment thickness and density data (on  $2 \times 2$  arc-deg grid) were used for computing the sediment (density contrast) stripping gravity correction with a spectral resolution complete to degree/order 90. CRUST2.0 (Bassin *et al.*, 2000) was compiled and administered by the US Geological Survey and the Institute for Geophysics and Planetary Physics at the University of California. CRUST2.0 is an upgrade of CRUST5.1 (Mooney *et al.*, 1998).

The Earth's solid topography (i.e. topographical heights onshore and bathymetric depths offshore) generated on a  $2 \times 2$  arc-deg grid using the DTM2006.0 coefficients complete to degree/order of 180 is shown in Figure 1. Maximum topographical heights over the study area reach  $4.1 \text{ km}$  and maximum bathymetric depths are  $5.7 \text{ km}$ .

Regional map of refined Bouguer gravity anomalies is shown in Figure 2, values were between  $-457$  and  $419 \text{ mGal}$  (mean =  $-43 \text{ mGal}$ ; standard deviation =  $286 \text{ mGal}$ ). Gravity maxima were found over areas of deepest oceans, and toand gravity minima were located throughout the central part of East Antarctica.

## Results

The combined least-squares method was used for a simultaneous estimation of Moho depth and density contrast. The solution was obtained by solving the system of normal equations, which was defined in equation (11). The observation vector  $\mathbf{I}$  in equation (10) was formed by three observation types  $l_1 = T\Delta\rho$  (eqn. 7),  $l_2 = \Delta\rho$  (eqn. 8), and  $l_3 = T_S$ . The initial values of Moho depths  $T_S$  were taken from

CRUST2.0 The variance-covariance matrix  $\mathbf{Q}$  in the least-squares estimation model (in eqn. 11) was computed from (cf. Sjöberg and Bagherabndi, 2011):

$$\mathbf{Q} = \begin{pmatrix} \sigma_1^2 & \sigma_1^2/t & 0 \\ \sigma_1^2/t & \sigma_2^2 & 0 \\ 0 & 0 & \sigma_3^2 \end{pmatrix}, \quad (14)$$

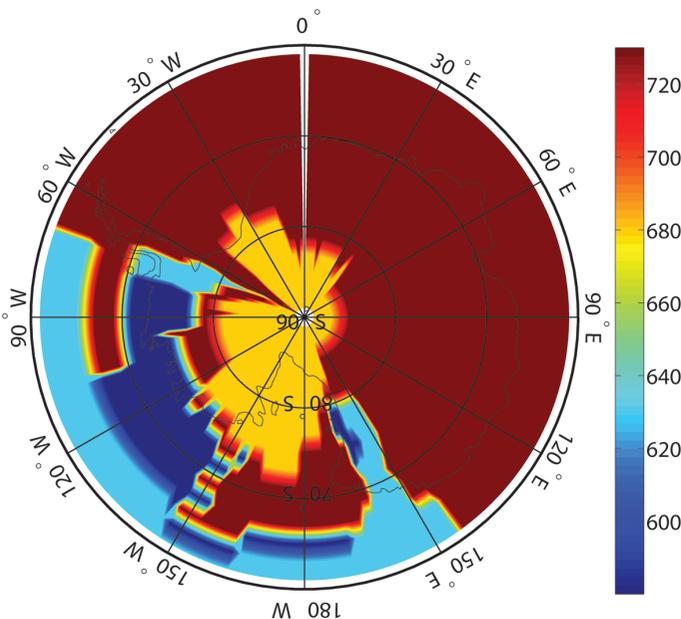
where  $\sigma_1$  and  $\sigma_3$  are the standard errors of  $T\Delta\rho$  and  $T$ , respectively, and  $\sigma_2^2 = \sigma_1^2/t^2 + \sigma_3^2(T\Delta\rho)^2/t^4$ . The standard error  $\sigma_1$  of  $T\Delta\rho$  was computed using the following expression:

$$\sigma_1^2 = \sigma_{T\Delta\rho}^2 \approx \left( \frac{\gamma_0}{4\pi G} \right)^2 \sum_{n,m} N_{n,m}^2 \sigma_{n,m}^2, \quad (15)$$

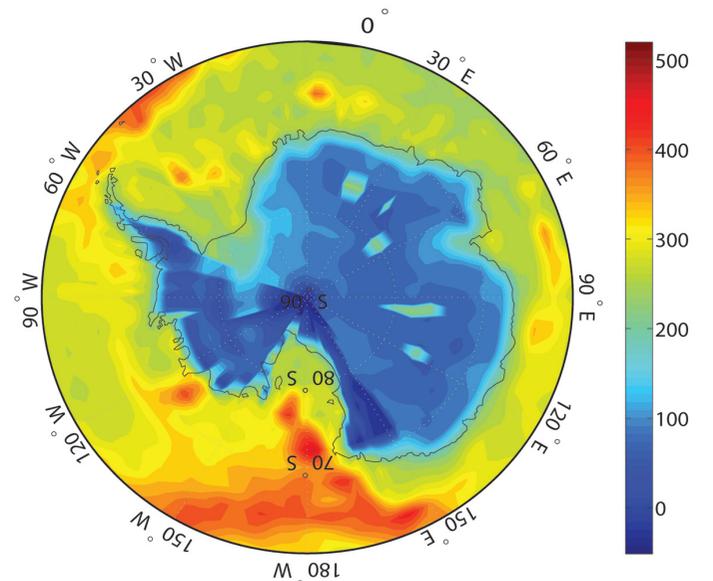
where  $\gamma_0$  is normal gravity,  $N_{n,m} = (2n+1)(n-1)/(n+1)$ , and  $\sigma_{n,m}^2$  error degree potential coefficients. Since CRUST2.0 Moho depths were not provided with a standard error model, relative uncertainties (i.e. standard errors  $\sigma_3$ ) of CRUST2.0 Moho depths of 20% were assumed in forming the matrix  $\mathbf{Q}$ .

Estimated Moho density contrast taken relative to  $2,670 \text{ kg/m}^3$  crust density are shown in Figure 3. The density contrast, was found to be between  $160$  and  $682 \text{ kg/m}^3$  (mean =  $477 \text{ kg/m}^3$ , standard deviation =  $128 \text{ kg/m}^3$ ).

Estimated Moho density contrast was compared with CRUST2.0. CRUST2.0 Moho density contrast was computed from CRUST2.0 upper mantle densities relative to  $2,670 \text{ kg/m}^3$  crust density. CRUST2.0 Moho density contrast is shown in Figure 4. Differences between both Moho density contrasts are shown in Figures 5. Whereas large variations of Moho density contrast were estimated using the combined least-squares approach (Fig. 2), the range of CRUST2.0 values was within  $580$  and  $730 \text{ kg/m}^3$  (Fig. 4). Such large discrepancies (shown in Fig. 5) were likely caused by a low quality of CRUST2.0 upper mantle density, especially under oceanic crust. CRUST2.0 upper mantle density under oceanic crust was significantly overestimated. CRUST2.0 Moho density contrast under the Antarctic continental crust had a good agreement with our estimates, differences mostly being within  $\pm 100 \text{ kg/m}^3$ .



**Fig. 4** CRUST2.0 Moho density contrast on a  $2 \times 2$  arc-deg grid. Values are in  $\text{kg/m}^3$ .



**Fig. 5** Differences between CRUST2.0 and newly estimated Moho density contrast on a  $2 \times 2$  arc-deg grid. Values are in  $\text{kg/m}^3$ .

## Discussion

Moho density contrast differs significantly in East and West Antarctica (see Fig. 3). Moho density contrast maxima are situated throughout the central part of East Antarctica with the extension under the Transantarctic mountain range. Moho density contrast there exceeded  $500 \text{ kg/m}^3$ . Moho density contrast beneath West Antarctica was less pronounced. A typical range was between  $400$  and  $500 \text{ kg/m}^3$ , except for local maxima in the central part of Antarctic Peninsula, where values reached  $\sim 550 \text{ kg/m}^3$ .

Ritzwoller *et al.*, (2001) and Block *et al.*, (2009) have estimated that average crust thickness in the West Antarctica was  $\sim 27$  km. Average crust thickness of  $\sim 40$  km was estimated for East Antarctica, with maximum crust thickness being  $\sim 45$  km. These large differences could be explained by different tectonic and geological compositions of West and East Antarctica. Dalziel and Elliot (1982) proposed that East Antarctica is a single tectonic block formed predominantly by Craton. Several smaller crust fragments forming West Antarctica have been moved and further reconfigured with respect to East Antarctica and each other. Ritzwoller *et al.*, (2001) and Morelli and Danesi (2004) have analysed shear wave velocities over the whole of Antarctica. They found significant differences in mantle structure characterised by a low velocity in the West Antarctica while a high velocity was detected in the mantle of East Antarctica, with a transition occurring below the Transantarctic Mountains. Unlike most of similar mountain ranges, this mountain range was formed in the absence of collision tectonic forces (ten Brink *et al.*, 1997). Several theories have been proposed explaining possible mechanisms for their formation. Models have included thermal buoyancy from an underlying positive temperature anomaly in upper mantle (ten Brink *et al.*, 1997), thicker crust giving the origin to an isostatically buoyant load (Studinger *et al.*, 2004) or possible collapse of a high-standing plateau with subsequent uplift and denudation (Bialas *et al.*, 2007). The latest models based on integrated geophysical analysis assume that multiple mechanisms have contributed to the uplift of the Transantarctic Mountains (cf. Studinger *et al.*, 2004; Lawrence *et al.*, 2006).

Despite the central part of East Antarctica being characterised by large variation in continental ice-sheet thickness and rough bedrock topography of basins and mountain ranges including sub-glacial mountains, significant variations in Moho density contrast were almost absent (see Fig. 3). Prevailing homogenous Moho density contrast there could be explained by a more uniform geological structure than that of West Antarctica. The location of the largest local density contrast anomalies correspond to the Gamburtsev subglacial mountains, Aurora Basin and Dronning Maud Land.

Block *et al.*, (2009) have reported decreasing crust thickness on the margins of East Antarctica due to rifting events that separated Antarctica from the other Gondwana landmasses, with the most pronounced signature occurring on the coast of Dronning Maud Land. This trend was not clearly seen in spatial pattern of Moho density contrast; instead, large Moho density contrast variations across the East Antarctic continental margin could be explained by significantly different geological composition between continental and oceanic crusts, involving typically thick and less dense continental crust compared to thin and heavier oceanic crust.

Distinctive features in the map of Moho density contrast geographically correspond to the West Antarctic Rift in West Antarctica; Moho density contrast there being typically  $400$  to  $450 \text{ kg/m}^3$ . Similar density contrast was found beneath the Ross Embayment. A possible explanation for these small values in Moho density contrast could be provided by geological structure dominated by volcanic activity occurring since (at least) the early Cenozoic. Holocene volcanism continues in the Ross Embayment (Kiele *et al.*, 1983; Blankenship *et al.*, 1993; Behrendt *et al.*, 1994). According to Behrendt *et al.*, (1994) the main rifting phase occurred between  $105$  and  $85$  Ma but episodic extension continued into the Cenozoic. Extension within the rift system has left most of West Antarctica below sea level, with the exception of Marie Byrd Land and parts of the Antarctic Peninsula. It is assumed that this represents remains of a continuously propagating rift

that started during the Jurassic period when Africa separated from East Antarctica and proceeded clockwise to its present location in the Ross Embayment and West Antarctica. The almost complete absence of recent seismic activity suggests that there might not be any undergoing active extension of the rift zone (Cande *et al.*, 2000). Local maxima of the Moho density contrast exceeding  $\sim 500 \text{ kg/m}^3$  in West Antarctica were found beneath Marie Byrd Land. The lithospheric structure there is characterised by topographic doming due to localised hot spot activity (Hole and LeMasurier, 1994; Winberry and Anandakrishnan, 2004).

Moho density contrast was typically less than  $\sim 350 \text{ kg/m}^3$  beneath oceanic crust; minima were found along the Pacific Antarctic mid-oceanic ridge. Moho density contrast there was estimated to be as low as  $160 \text{ kg/m}^3$ , corresponding to  $2,830 \text{ kg/m}^3$  density for the youngest oceanic lithosphere (Müller *et al.*, 2008). Similar density values have also been found between Antarctic and South-American continental tectonic plates, which were separated by a newly-formed oceanic lithosphere.

## Conclusions

Spatial variations of Moho density contrast resembled major features of the Antarctic tectonic plate configuration and its geological composition. Density contrast maxima were found throughout the central part of East Antarctica with the extension under the Transantarctic mountain range. Moho density contrast there often exceeded  $\sim 500 \text{ kg/m}^3$ , maxima being  $682 \text{ kg/m}^3$ . Density contrast maxima were found under mountainous ranges and areas covered by the largest continental ice sheet. These locations are also characterised by the largest thickness of the Antarctic continental crust, Moho depths there reached  $\sim 40$  km.

According to our estimation, Moho density contrast beneath West Antarctica was typically  $400$  to  $500 \text{ kg/m}^3$ , except for local maxima ( $\sim 550 \text{ kg/m}^3$ ) in the central part of Antarctic Peninsula. Local minima ( $400$ – $450 \text{ kg/m}^3$ ) beneath the West Antarctic Rift and the Ross Embayment are likely attributed to volcanic compositions throughout this divergent tectonic zone.

The largest variations of Moho density contrast were found across the East Antarctic continental margin. Moho density contrast beneath the oceanic crust was typically less than  $\sim 350 \text{ kg/m}^3$ . The extreme minima were detected along the Pacific Antarctic mid-oceanic ridge.

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