

# Green's Functions for Sturm-Liouville Problems on Directed Tree Graphs

Funciones de Green para problemas de Sturm-Liouville en árboles direccionales

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**ABSTRACT.** Let  $\Gamma$  be geometric tree graph with  $m$  edges and consider the second order Sturm-Liouville operator  $\mathcal{L}[u] = (-pu')' + qu$  acting on functions that are continuous on all of  $\Gamma$ , and twice continuously differentiable in the interior of each edge. The functions  $p$  and  $q$  are assumed continuous on each edge, and  $p$  strictly positive on  $\Gamma$ . The problem is to find a solution  $f : \Gamma \rightarrow \mathbb{R}$  to the problem  $\mathcal{L}[f] = h$  with  $2m$  additional conditions at the nodes of  $\Gamma$ . These node conditions include continuity at internal nodes, and jump conditions on the derivatives of  $f$  with respect to a positive measure  $\rho$ . Node conditions are given in the form of linear functionals  $l_1, \dots, l_{2m}$  acting on the space of admissible functions. A novel formula is given for the Green's function  $G : \Gamma \times \Gamma \rightarrow \mathbb{R}$  associated to this problem. Namely, the solution to the semi-homogenous problem  $\mathcal{L}[f] = h$ ,  $l_i[f] = 0$  for  $i = 1, \dots, 2m$  is given by  $f(x) = \int_{\Gamma} G(x, y)h(y) d\rho$ .

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**RESUMEN.** Sea  $\Gamma$  un grafo tipo árbol con  $m$  aristas y considere el operador de Sturm-Liouville  $\mathcal{L}[u] = (-pu')' + qu$  definido en el espacio de funciones continuas en  $\Gamma$  y continuamente diferenciables dos veces al interior de cada arista de  $\Gamma$ . Las funciones  $p$  y  $q$  se suponen continuas en cada arista, y  $p$  es estrictamente positiva en todo  $\Gamma$ . El problema consiste en hallar la solución  $f : \Gamma \rightarrow \mathbb{R}$  al problema dado por  $\mathcal{L}[f] = h$  mas  $2m$  condiciones en los nodos de  $\Gamma$ : en los nodos internos se especifican continuidad de  $f$  y condiciones de salto para las derivadas de  $f$  con respecto a una medida  $\rho$ . Estas condiciones de nodo se expresan en la forma de funcionales lineales  $l_1, \dots, l_{2m}$  actuando sobre el espacio de funciones admisibles para  $\mathcal{L}$ . Se presenta una nueva fórmula

para la función de Green  $G : \Gamma \times \Gamma \rightarrow \mathbb{R}$  asociada con este problema. Es decir, se expresa la solución del problema semi-homogéneo  $\mathcal{L}[f] = h$ ,  $l_i[f] = 0$  para  $i = 1, \dots, 2m$  como  $f(x) = \int_{\Gamma} G(x, y)h(y) d\rho$ .

*Palabras y frases clave.* Sturm-Liouville problems on graphs, Green's function.

## 1. Introduction

The Sturm-Liouville differential operator

$$\mathcal{L}[f] = -(pf')' + qf \quad (1)$$

on an interval, appears in the analysis of many different types of models in the natural sciences. The problem  $\mathcal{L}[f] = h$  or  $\mathcal{L}[f] = \nu qf$ , together with appropriate boundary conditions, arises when considering Kirchoff's law in electrical circuits, the balance of tension in a elastic string, or the steady state temperature in a heated rod (see for example [4, 3, 2]). A more complete review of the mathematical theory can be found in [10].

The extension of operator (1) to the case of a domain composed of intervals arranged in a graph has received recent attention (see for example [5, 9, 1]). A complete bibliographical review with historical notes and applications can be found in [6]. The particular work presented here is motivated by the following problem in mathematical ecology: the stability of populations of organisms in river networks where the dispersion of individuals is governed by an advection-diffusion operator.

is motivated by an applied problem in mathematical ecology, that of the stability for populations of organisms in river networks [8].

Let  $\Gamma$  be a tree graph in  $\mathbb{R}^2$ , that is,  $\Gamma$  is a collection of edges joined with nodes, where each edge can be treated as a finite open interval (a complete description of the notation and assumptions for  $\Gamma$  is laid out in Section 2.1). We are interested in the Storm-Liouville operator of the form

$$\mathcal{L}[f]_e = -(pf'_e)' + qf_e, \quad e \in \Gamma, \quad (2)$$

where the subscript  $e$  denotes the restriction of a function to edge  $e$ . In Equation (2), the functions  $p$  and  $q$  are assumed continuous and bounded, with  $p$  being also uniformly bounded away from zero.

Given a function  $h : \Gamma \rightarrow \mathbb{R}$ , one is interested in solving the problem

$$\mathcal{L}[f] = h, \quad f \in \mathcal{E}(\Gamma) \quad (3)$$

where  $\mathcal{E}(\Gamma)$  denotes the set of functions that are twice continuously differentiable inside each edge, and satisfy a given set of node conditions. It can easily be shown that problem (3) has a unique solution if and only if it is non-degenerate, namely, the only solution to the homogeneous problem  $\mathcal{L}[f] = 0$ ,  $f \in \mathcal{E}(\Gamma)$ , is  $f \equiv 0$ .

Let  $\mathcal{G}$  be some right inverse mapping of  $\mathcal{L}$  such that  $\mathcal{L}[\mathcal{G}[f]] = f$  for all admissible  $f$ . Then  $\mathcal{G}$  is the Green's operator for problem (3). Moreover, it will be shown that:

**Theorem 1.** *If problem (3) is non-degenerate, it has a unique Green's function. Namely, there exists a function  $G : \Gamma \times \Gamma \rightarrow \mathbb{R}$  such that the solution to (3),  $f = \mathcal{G}[h]$ , is  $f(x) = \int_{\Gamma} G(x, y)h(y) dy$ .*

The proof of the existence of the Green's function  $G$  defined in Theorem 1 follows from standard arguments. Here, the proof is obtained by simply giving an explicit formula for  $G(x, y)$ . Uniqueness of the Green's function follows from the hypothesis of non-degeneracy.

In [7] the authors provide a proof of existence, uniqueness and a formula for  $G$  for general graphs. The goal here is to present a new formula, both simpler and less expensive to compute, for the Green's function in the case  $\Gamma$  is a tree graph. The techniques here are elementary and based on the classical Lagrange's method for Sturm-Liouville problems (see for example [2]).

The organization is as follows. The next Section settles the notation and defines the class of Sturm-Liouville problems to be considered. Finally, Section 3 is devoted to the construction and the formula for the the Green's function.

## 2. Sturm-Liouville Problems on Tree Graphs

### 2.1. Tree Graphs and Functions

By a tree graph we understand a finite collection of edges embedded in  $\mathbb{R}^2$ , joined with nodes and containing no loops. That is, for any two points  $x, y$  in the graph, there exists one single path through the graph joining them. We assume that each edge  $e$  of the graph allows a sufficiently smooth parametrization, contains no self-intersections, and is finite, therefore can be considered as the interval  $e = (0, l_e)$ . The collection of all edges is denoted by  $\Gamma$ . At each endpoint of an edge it is located a node of  $\Gamma$ . The set of nodes is  $N(\Gamma)$  and boldface is used to denote individual nodes. The graph, including its nodes, is denoted as  $\bar{\Gamma} := \Gamma \cup N(\Gamma)$ .

Points in  $\Gamma$  are denoted by the pair  $(e, x)$  with  $0 < x < l_e$ , or by single letters if specification of the edge is not crucial. If  $\mathbf{n}$  is a node, let  $i(\mathbf{n})$  denote the set of incident edges at  $\mathbf{n}$ , namely those for which  $\mathbf{n}$  is an endpoint. Boundary nodes are those  $\mathbf{n}$  with  $\#i(\mathbf{n}) = 1$ . The set of all boundary nodes of  $\Gamma$  is  $\partial\Gamma$ . The set of internal nodes is  $I(\Gamma) = N(\Gamma) \setminus \partial\Gamma$ . Node  $\mathbf{n}$  has  $\#i(\mathbf{n}) \geq 1$  possible representations: for each  $e \in i(\mathbf{n})$ , one either has  $\mathbf{n} = (e, 0)$  or  $\mathbf{n} = (e, l_e)$ . The representation of points in  $\bar{\Gamma}$  is therefore dependent on the parametrization direction of its edges.

The value of a function  $f : \Gamma \rightarrow \mathbb{R}$  at a point in  $\Gamma$  is denoted as  $f_e(x) = f(e, x)$ . That is,  $f_e$  is the restriction of  $f$  to the edge  $e$ . For a node  $\mathbf{n} \in \partial\Gamma$

located at the endpoint of some edge  $e$ ,  $f(\mathbf{n})$  denotes the appropriate one-sided limit of  $f_e$ . For an internal node  $\mathbf{n}$  with  $i(\mathbf{n}) = \{e_1, \dots, e_n\}$  the value  $f_{e_1}(\mathbf{n})$  denotes the one-sided limit of  $f_{e_i}$  as  $x$  approaches the endpoint of  $e_i$  at which  $\mathbf{n}$  is located,  $i = 1, \dots, n$ . If all these limits coincide,  $f$  is said to be continuous at  $\mathbf{n}$ , and  $f(\mathbf{n})$  is defined as the common value.

We must also differentiate functions given on  $\Gamma$ . For a point  $(e, x) \in \Gamma$  with  $0 < x < l_e$ , the derivative  $f'(e, x) = f'_e(x)$  is computed as the usual derivative of the restriction  $f_e$  at  $x$  according to the particular parametrization direction of  $e$ . A change in the orientation of the parametrization of the edge implies a sign change on  $f'_e$ . Note that the sign of  $f''_e$  or  $(pf'_e)'$  remains unchanged. For a node  $\mathbf{n}$  located at an endpoint of edge  $e$ , we introduce the *boundary derivative*  $f_e^{fb}$  as the derivative “out of node  $\mathbf{n}$  into edge  $e$ ”, as if the parametrization of  $e$  has  $\mathbf{n} = (e, 0)$ . Boundary derivatives are useful because they make the following equality hold

$$\int_e (pf')' dx = p(l_e)f^{fb}(l_e) - p(0)f^{fb}(0), \quad (4)$$

regardless of whether the integral is computed from 0 to  $l_e$ , or from  $l_e$  to 0.

The space of functions that are  $n$  times continuously differentiable in  $\Gamma$  is denoted by  $\mathcal{C}^n(\Gamma)$ ,  $n = 0, 1, 2, \dots$ ;  $\mathcal{C}(\Gamma) := \mathcal{C}^0(\Gamma)$ . Clearly, such spaces are identifiable with direct sums of the form  $\bigoplus_e \mathcal{C}^n(0, l_e)$ . The set  $\mathcal{C}(\bar{\Gamma})$  is composed of functions in  $\mathcal{C}(\Gamma)$  that are also continuous at each node.

## 2.2. Sturm-Liouville Operators

Let  $p, q \in \mathcal{C}(\Gamma)$  be bounded with  $\inf_{\Gamma} p(x) > 0$ . The object of this study is the differential operator

$$\mathcal{L}[f] := -(pf')' + qf, \quad f \in \mathcal{D}_p^2(\Gamma), \quad (5)$$

where  $\mathcal{D}_p^2(\Gamma)$  denotes the space of functions  $f \in \mathcal{C}(\Gamma)$  such that  $(pf')' \in \mathcal{C}(\Gamma)$ .

If  $\Gamma$  contains  $m$  edges, then the dimension of  $\mathcal{N}(\mathcal{L}) := \{u \in \mathcal{D}_p^2(\Gamma) : \mathcal{L}[u] = 0\}$  is  $2m$ . A basis  $\{\varphi_1, \dots, \varphi_{2m}\}$  for  $\mathcal{N}(\mathcal{L})$  can be found as follows. Let  $e$  be the  $i$ -th edge and define  $\varphi_{2i-1}$  and  $\varphi_{2i}$  as the solutions to  $-(p_e f')' + q_e f = 0$  on  $e$  satisfying

$$\varphi_{2i-1}(0) = 1, \quad \varphi'_{2i-1}(0) = 0, \quad \varphi_{2i}(0) = 0, \quad \varphi'_{2i}(0) = 1,$$

extended to all of  $\Gamma$  via  $\varphi_{2i-1}(\tilde{e}, x) = \varphi_{2i}(\tilde{e}, x) = 0$ ,  $0 < x < l_{\tilde{e}}$  for all  $\tilde{e} \neq e$ . Consider now a collection of  $2m$  linear functionals  $\{l_i, \dots, l_{2m}\}$  defined on  $\mathcal{D}_p^2(\Gamma)$ . These functionals play the role of the linear conditions at the nodes of  $\Gamma$ . The problem

$$f \in \mathcal{D}_p^2(\Gamma), \quad \mathcal{L}[f] = h, \quad l_i[f] = c_i, \quad i = 1, \dots, 2m \quad (6)$$

will be uniquely solvable if and only if the homogenous problem

$$f \in \mathcal{D}_p^2(\Gamma), \quad \mathcal{L}[f] = 0, \quad l_i[f] = 0, \quad i = 1, \dots, 2m \quad (7)$$

has no solution except the trivial solution  $f \equiv 0$ . In this case we say that problem (6) is *non-degenerate*.

Non-degeneracy can be characterized as follows. Let  $\Delta$  be the matrix defined by  $\Delta_{i,j} = l_i[\varphi_j]$ ,  $i, j = 1, \dots, 2m$ . Non-degeneracy is therefore equivalent to  $\det(\Delta) \neq 0$ . In this case, the solution to problem (6) can be written explicitly. Let  $z$  be some solution to the semi-homogeneous problem

$$z \in \mathcal{D}_p^2(\Gamma), \quad \mathcal{L}[z] = h, \quad l_i[z] = 0, \quad i = 1, \dots, 2m. \quad (8)$$

Then, the solution  $f = z + \sum_{i=1}^{2m} a_i \varphi_i$  to problem (6) must satisfy

$$\begin{bmatrix} 1 & \varphi_1 & \cdots & \varphi_{2m} \\ 0 & \boxed{\Delta} \\ \vdots & & & \\ 0 & & & \end{bmatrix} \begin{bmatrix} f \\ -a_1 \\ \vdots \\ -a_{2m} \end{bmatrix} = \begin{bmatrix} z \\ -c_1 \\ \vdots \\ -c_{2m} \end{bmatrix}. \quad (9)$$

Hence, Cramer's rule gives the useful formula

$$f = \frac{1}{\det(\Delta)} \det \begin{bmatrix} z & \varphi_1 & \cdots & \varphi_{2m} \\ -c_1 & \boxed{\Delta} \\ \vdots & & & \\ -c_{2m} & & & \end{bmatrix}. \quad (10)$$

### 2.3. The Physical Problem

Motivated by physical applications, we now specialize to semi-homogenous Sturm-Liouville problems where the collections of functionals  $\{l_i : i = 1, \dots, 2m\}$  corresponds to a particular choice of conditions at the nodes of  $\bar{\Gamma}$ .

Consider the operator  $\mathcal{L}[f]$  acting on the set

$$\mathcal{E}(\Gamma) = \mathcal{D}_p^2(\Gamma) \cap \mathcal{C}(\bar{\Gamma}) \cap F_\rho(\Gamma) \cap B_D(\Gamma), \quad (11)$$

where  $B_D(\Gamma)$  and  $F_\rho(\Gamma)$  contain the boundary and weighted flux matching conditions respectively:

$$B_D = \{f \in \mathcal{C}(\bar{\Gamma}) : f(\mathbf{n}) = 0, \mathbf{n} \in \partial\Gamma\}, \quad (12)$$

$$F_\rho(\Gamma) = \left\{ f \in \mathcal{C}^1(\Gamma) : \sum_{e \in (\mathbf{n})} \rho_e f_e'(\mathbf{n}) = 0, \mathbf{n} \in I(\Gamma) \right\}. \quad (13)$$

The function  $\rho$  in (12) is assumed constant on edges and strictly positive. Other linear boundary conditions types than Dirichlet can be considered without major changes to the arguments that follow.

The conditions encoded in  $\mathcal{E}(\Gamma)$  can be cast in terms of linear functionals: let  $\mathbf{n} \in I(\Gamma)$  with  $i(\mathbf{n}) = \{e_1, \dots, e_k\}$ , and define the functionals

$$\tilde{l}_{\mathbf{n},i}[f] = f_{e_{i+1}}(\mathbf{n}) - f_{e_i}(\mathbf{n}), \quad i = 1, \dots, k-1, \quad (14)$$

$$\tilde{l}_{\mathbf{n},k}[f] = \sum_{i=1}^k \rho_{e_i} f'_{e_i}(\mathbf{n}). \quad (15)$$

For a boundary node  $\mathbf{n}$  located at the endpoint of edge  $e$ , define simply

$$\tilde{l}_{\mathbf{n}}[f] = f_e(\mathbf{n}). \quad (16)$$

Relabeling gives a collection of  $2m$  functionals  $\{l_1, \dots, l_{2m}\}$ , such that the problem of finding  $f \in \mathcal{E}(\Gamma)$  satisfying  $\mathcal{L}[f] = h$  can be written as the semi-homogenous problem (8).

### 3. Construction of the Green's Function

The goal is to arrive to a formula for the solution to problem (8). The solution to the associated non-homogenous problem will then follow from (10).

**Definition 2.** A Green's function for operator  $\mathcal{L}$  is a function  $G : \Gamma \times \Gamma \rightarrow \mathbb{R}$  such that for all  $h \in \text{Ran}(\mathcal{L})$ , the solution to problem (8) is given by

$$f(x) = \int_{\Gamma} G(x, y) h(y) dy. \quad (17)$$

The operator  $\mathcal{G} : h \mapsto \int_{\Gamma} G(x, y) h(y) dy$  is called the Green's operator.

The first step is elementary and consists on verifying properties of the Wronskian of functions on  $\Gamma$ . For  $f, g \in \mathcal{C}^1(\Gamma)$  the Wronskian  $W[f, g]$  is defined on an edge  $e$  of  $\Gamma$  as

$$W[f, g]_e = f_e g'_e - g_e f'_e. \quad (18)$$

**Lemma 3.** Let  $f, g, h$  be functions in  $\mathcal{C}^1(\Gamma)$ .

a) If  $f, g \in \mathcal{D}_p^2(\Gamma)$  then Lagrange's identity holds on each edge  $e$ ,

$$f_e \mathcal{L}[g]_e - g_e \mathcal{L}[f]_e = - \frac{d}{dx} (p_e W[f, g]_e). \quad (19)$$

b) If  $f, g \in B_D$ , then  $W[f, g] \in B_D(\Gamma)$ .

- c) If  $f, g \in \mathcal{C}(\bar{\Gamma}) \cap F_\rho(\Gamma)$ , then  $\sum_{e \in i(\mathbf{n})} \rho_e W[f, g]_e(\mathbf{n}) = 0$  for all  $\mathbf{n} \in I(\Gamma)$ . Here, the derivatives in the definition  $W$  at  $\mathbf{n}$  are replaced by boundary derivatives.
- d)  $hW[f, g] - fW[h, g] = gW[f, h]$ .
- e) If  $f, g \in \mathcal{D}_p^2(\Gamma)$  with  $\mathcal{L}[f] = \mathcal{L}[g] = 0$  on some edge, then  $pW[f, g]$  is constant there.

**Proof.** Statement a) follows from a simple calculation and b) is obvious. For c), it suffices to use continuity and change derivatives to boundary derivatives,

$$\sum_{e \in i(\mathbf{n})} \rho_e W[f_e, g_e](\mathbf{n}) = f(\mathbf{n}) \sum_{e \in i(\mathbf{n})} \rho_e g_e'^b(\mathbf{n}) + g(\mathbf{n}) \sum_{e \in i(\mathbf{n})} \rho_e f_e'^b(\mathbf{n}) = 0.$$

d) is obtained by rearranging terms. To prove e), compute  $(p_e W[f, g]_e)' = f_e(p_e g_e')' - g_e(p_e f_e')'$ , use  $(p_e f_e')' = -q_e f_e$  and  $(p_e g_e')' = -q_e g_e$  to finally get  $(p_e W[f, g]_e)' = 0$ .  $\square$

For the following definition, and subsequent formulas, assume without loss of generality that the parametrization of  $\Gamma$  is such that for all  $\mathbf{n} \in \partial\Gamma$ , we have  $\mathbf{n} = (e, l_e)$ , where  $e$  is the edge to which  $\mathbf{n}$  belongs.

**Definition 4.** Refer to Figure 1. Let  $(e, x) \in \Gamma$ ,

- (1) The two connected components of  $\bar{\Gamma} \setminus \{(e, x)\}$  are denoted  $\bar{\Gamma}(e, x)$  and  $\bar{\Lambda}(e, x)$  respectively. The point  $(e, x)$  is adjoined as a boundary node to  $\bar{\Gamma}(e, x)$  and  $\bar{\Lambda}(e, x)$ . By convention,  $\bar{\Gamma}(e, x)$  is taken as the tree that contains the node  $(e, 0)$ . There is an edge denoted by  $e$  in both  $\bar{\Gamma}(e, x)$  and  $\bar{\Lambda}(e, x)$ ; it is parametrized as the intervals  $(0, x)$  and  $(x, l_e)$  respectively.
- (2) For an edge  $e$ ,  $\bar{\Gamma}(e) := \cup_{x \in (0, l_e)} \bar{\Gamma}(e, x)$ , and  $\bar{\Lambda}(e) := \cup_{x \in (0, l_e)} \bar{\Lambda}(e, x)$ .
- (3) As in the case of the full tree,  $\Gamma(e, x)$  – without the bar – denotes the collection of points inside edges of  $\bar{\Gamma}(e, x)$ . Similarly for  $\Lambda(e, x)$ ,  $\Gamma(e)$ ,  $\Lambda(e)$ .
- (4)  $\mathcal{E}_0(\Gamma(e))$  is the set of functions  $f \in \mathcal{D}_p^2(\Gamma(e)) \cap \mathcal{C}(\bar{\Gamma}(e)) \cap F_\rho(\Gamma(e))$  such that  $f(\mathbf{n}) = 0$  for  $\mathbf{n} \in \partial\Gamma(e) \setminus \{(e, l_e)\}$ .
- (5) Similarly,  $\mathcal{E}_0(\Lambda(e))$  is comprised of functions  $f \in \mathcal{D}_p^2(\Lambda(e)) \cap \mathcal{C}(\bar{\Lambda}(e)) \cap F_\rho(\Lambda(e))$  such that  $f(\mathbf{n}) = 0$  for  $\mathbf{n} \in \partial\Lambda(e) \setminus \{(e, 0)\}$ .

**Lemma 5.** Let  $e$  be a fixed edge. If problem (8) is non-degenerate, then there exist solutions  $\psi^{\Gamma(e)} \in \mathcal{E}_0(\Gamma(e))$ ,  $\psi^{\Lambda(e)} \in \mathcal{E}_0(\Lambda(e))$  to  $\mathcal{L}\psi^{\Gamma(e)} = 0$  on  $\Gamma(e)$ , and  $\mathcal{L}\psi^{\Lambda(e)} = 0$  on  $\Lambda(e)$ . These functions can further be chosen such that  $pW[\psi^{\Gamma(e)}, \psi^{\Lambda(e)}] = -1$  on  $e$ .

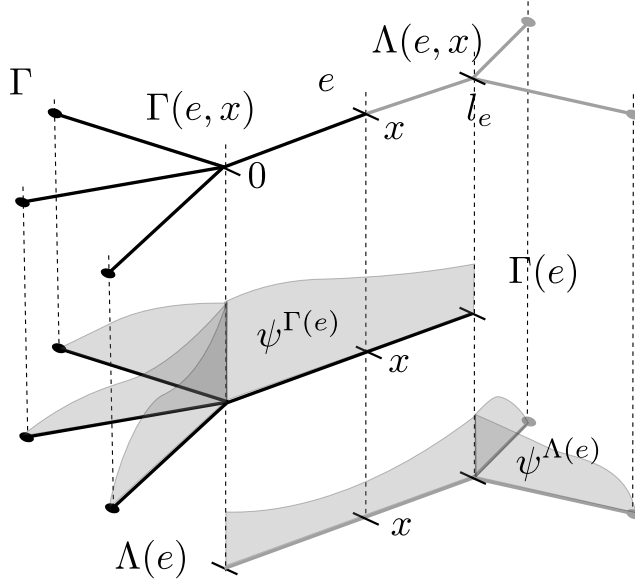


FIGURE 1. Schematic representation of a tree  $\Gamma$  with six edges. For  $e$  and  $x$  as shown, the black sub-tree on the upper figure in  $\Gamma(e, x)$ , the gray sub-tree is  $\Lambda(e, x)$ . The middle and lower figures depict trees  $\Gamma(e)$  and  $\Lambda(e)$  respectively with a schematization of the functions  $\psi^{\Gamma(e)}$  and  $\psi^{\Lambda(e)}$ .

**Proof.** If  $e$  contains no nodes in  $\partial\Gamma$ , choose a node  $\mathbf{n}' \notin \partial\Gamma(e)$  and let  $e'$  be its edge. If  $e$  contains a node in  $\partial\Gamma$ , make  $\mathbf{n}'$  equal to that node, and  $e' = e$ . Rearrange the basis  $\{\varphi_1, \dots, \varphi_{2m}\}$  such that  $\varphi_1$  and  $\varphi_2$  are supported on  $e'$ , and the collection of functionals  $\{l_1, \dots, l_{2m}\}$  such that  $l_1[f] = f(\mathbf{n}')$ . Let  $\Delta$  be the matrix defined in Section (2.2). By the nondegeneracy of problem (8), there exists a solution  $a = (a_1, \dots, a_{2m})$  to  $\Delta a = \varepsilon^{(1)}$ , where  $\varepsilon^{(1)}$  denotes the  $\mathbb{R}^{2m}$  vector that has one in the first coordinate, and zero elsewhere. The function  $\psi = \sum_{i=1}^{2m} a_i \varphi_i$  is a solution in  $\mathcal{D}_p^2(\Gamma) \cap \mathcal{C}(\bar{\Gamma}) \cap F_p(\Gamma)$  to  $\mathcal{L}[\psi] = 0$  on all of  $\Gamma$  and such that  $\psi(\mathbf{n}) = 0$  for  $\mathbf{n} \in \partial(\Gamma) \setminus \{\mathbf{n}'\}$ . The restriction of  $\psi$  to  $\Gamma(e)$  serves as the required function  $\psi^{\Gamma(e)}$ . A similar construction applies for  $\psi^{\Lambda(e)}$ . By Lemma 3,  $pW[\psi^{\Gamma(e)}, \psi^{\Lambda(e)}]$  is constant on  $e$ , and the desired normalization can be achieved if this constant is not zero. Assume on the contrary that  $pW[\psi^{\Gamma(e)}, \psi^{\Lambda(e)}] = 0$  on  $e$ . Since the Wronskian vanishes, there is  $k \neq 0$  such that  $\psi_e^{\Gamma(e)} = k\psi_e^{\Lambda(e)}$ . Then, the function  $f := \psi^{\Gamma(e)}\mathbf{1}_{\Gamma(e)} + k\psi^{\Lambda(e)}\mathbf{1}_{\Gamma(e)^c}$  would be a solution to the homogenous problem (7) violating the assumption of non-degeneracy.  $\square$



**Remark 6.** The computation of the  $\psi^{\Lambda(e)}, \psi^{\Gamma(e)}$  can be performed quite inexpensively. For a boundary node  $\mathbf{n}' = (e', l_{e'}) \in \partial\Gamma$ , the solution  $\psi$  constructed in the proof of Lemma 5 can be restricted to define  $\psi^{\Gamma(e)}$  for all nodes  $e$  such that either  $e = e'$  or  $e'$  does not belong to  $\Gamma(e)$ . Similarly, it can be used to define  $\psi^{\Lambda(e)}$  for all nodes  $e$  such that  $e'$  does not belong to  $\Lambda(e)$ . This implies that the linear system  $\Delta a = \varepsilon^{(1)}$  has to be solved only  $\#\partial\Gamma$  times.

The specific form of the Green's function can now be written.

**Theorem 7.** *Assume problem (8) is non-degenerate. The following function is a Green's function for operator  $\mathcal{L}$*

$$G(x, y) = \frac{1}{\rho_e} \times \begin{cases} \psi^{\Gamma(e)}(y) \psi^{\Lambda(e)}(x), & y \in \Gamma(e, x) \\ \psi^{\Lambda(e)}(y) \psi^{\Gamma(e)}(x), & y \in \Lambda(e, x) \end{cases}, \quad x \in e. \quad (20)$$

Moreover, this function is unique in the class of continuous functions on  $\Gamma$  that are continuous with respect to the first variable.

**Proof.** Let  $h \in \text{Ran}(\mathcal{L})$ , and  $f \in \mathcal{E}(\Gamma)$  a solution to  $\mathcal{L}[f] = h$ . Fix an edge  $e$ , and  $x \in e$ . Applying Lagrange's identity (19) for  $\psi^{\Gamma(e)}$  and  $f$  and integrating over  $\Gamma(e, x)$  with respect to the measure  $\rho$  gives

$$\int_{\Gamma(e, x)} \psi^{\Gamma(e)} h \, d\rho = - \sum_{a \subset \Gamma(e, x)} \left( p_a W[\psi^{\Gamma(e)}, f]_a \rho_a \right) \Big|_0^{l_a},$$

where the sum on the right hand side is taken over all edges  $a$  of  $\Gamma(e, x)$ . Parts b) and c) of Lemma 3 ensure that all terms in the sum cancel except for the value at  $(e, x)$ ,

$$\int_{\Gamma(e, x)} \psi^{\Gamma(e)} h \, d\rho = -p_e(x) \rho_e W[\psi^{\Gamma(e)}, f]_e(x). \quad (21)$$

Similarly, Lagrange's identity for  $\psi^{\Lambda(e)}$  and  $f$ , gives

$$\int_{\Lambda(e, x)} \psi^{\Lambda(e)} h \, d\rho = p_e(x) \rho_e W[\psi^{\Lambda(e)}, f]_e(x). \quad (22)$$

Multiply equations (21) and (22) by  $\psi^{\Lambda(e)}(x)$  and  $\psi^{\Gamma(e)}(x)$  respectively, add the resulting equations, and apply part d) of Lemma 3 to the right hand side of the result. Finally, since  $pW[\psi^{\Gamma(e)}, \psi^{\Lambda(e)}] = -1$  on  $e$ ,

$$\int_{\Gamma(e,x)} \psi^{\Lambda(e)}(x) \psi^{\Gamma(e)}(y) h(y) d\rho(y) + \int_{\Lambda(e,x)} \psi^{\Gamma(e)}(x) \psi^{\Lambda(e)}(y) h(y) d\rho(y) = f_e(x) \rho_e. \quad (23)$$

Since  $\Gamma$  is a disjoint union of  $\Gamma(e, x)$  and  $\Lambda(e, x)$ , the function  $G(x, y)$  defined in (20) satisfies Definition (2). Let  $h \in C(\Gamma)$  be arbitrary. It will be established now that  $h \in \text{Ran}(\mathcal{L})$  simply by showing that  $f := \mathcal{G}h$  solves  $\mathcal{L}[f] = h$ . Write

$$f(x) = \psi^{\Lambda(e)}(x) \int_{\Gamma(e) \setminus e} \psi^{\Gamma(e)} h d\rho + \psi^{\Gamma(e)}(x) \int_{\Lambda(e) \setminus e} \psi^{\Lambda(e)} h d\rho + \psi^{\Lambda(e)}(x) \int_0^x \psi^{\Gamma(e)} h d\rho + \psi^{\Gamma(e)}(x) \int_x^{l_e} \psi^{\Lambda(e)} h d\rho. \quad (24)$$

Applying  $\mathcal{L}$  to the first two terms in (24) gives zero since  $\mathcal{L}[\psi^{\Lambda(e)}] = \mathcal{L}[\psi^{\Gamma(e)}] = 0$ . A routine calculation finally shows that

$$\mathcal{L}[f] = -hpW[\psi^{\Gamma(e)}, \psi^{\Lambda(e)}]_e + \mathcal{L}[\psi^{\Lambda(e)}] \int_0^x \psi^{\Gamma(e)} h d\rho + \mathcal{L}[\psi^{\Gamma(e)}] \int_x^{l_e} \psi^{\Lambda(e)} h d\rho$$

which yields  $\mathcal{L}[f] = h$ . Lastly, the non-degeneracy of problem (8) and the fact that  $\text{Ran}(\mathcal{L}) = \mathcal{C}(\Gamma)$ , imply the uniqueness of  $G$  as stated in the theorem.  $\square$

**Remark 8.** The construction of the Green's function in Theorem 7 has one particular important advantage over the one proposed by [7]. In that work,  $G(x, y)$  is given as

$$G(x, y) = H(x, y) - \sum_{i=1}^{2m} l_i [H(\cdot, y)] \eta_i(x) \quad (25)$$

where  $H(x, y)$  is equal to the Green's function of operator  $\mathcal{L}$  on  $(0, l_e)$  if  $x, y \in e$ , and equal to zero whenever  $x$  and  $y$  belong to different edges. The functions  $\eta_i$  are solutions to  $\mathcal{L}[\eta_i] = 0$ ,  $l_j[\eta_i] = \delta_{ij}$ . Note that this formula requires solving  $\Delta a = \varepsilon^{(1)}$  a total of  $2m$  times to compute  $G(x, y)$  at any given single pair of points  $(e_x, x)$ ,  $(e_y, y)$  of  $\Gamma$ . Via formula (20), one needs only the functions  $\psi^{\Lambda(e_x)}$  and  $\psi^{\Gamma(e_x)}$  and therefore, the system  $\Delta a = \varepsilon^{(1)}$  must be solved only twice. On the other hand, formula (25) has the advantage of using  $H$ , which is a diagonal fundamental solution to  $\mathcal{L}[f] = h$ .

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